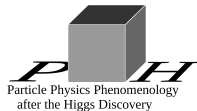


New Physics in $b \rightarrow c\tau\nu$: Impact of Polarisation Observables and $B_c \rightarrow \tau\nu$

Monika Blanke

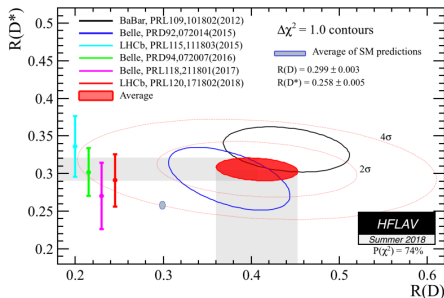


LHCb and Belle II Opportunities for Model Builders
MITP – January 31, 2019

The $\mathcal{R}(D^{(*)})$ Anomaly

Test of lepton flavour universality in semi-leptonic B decays

$$\mathcal{R}(D^{(*)}) = \frac{\text{BR}(B \rightarrow D^{(*)} \tau \nu)}{\text{BR}(B \rightarrow D^{(*)} \ell \nu)} \quad (\ell = e, \mu)$$



- **theoretically clean**, as hadronic uncertainties largely cancel in ratio
- measurements by BaBar, Belle, and LHCb (so far $\mathcal{R}(D^*)$ only)
- **3.8 σ tension** between HFLAV fit and SM value
- (qualitatively) supported by measurement of $\mathcal{R}(J/\psi)$ (LHCb)

Related Observables

- ratio of baryonic decay rates

$$\mathcal{R}(\Lambda_c) = \frac{\text{BR}(\Lambda_b \rightarrow \Lambda_c \tau \nu)}{\text{BR}(\Lambda_b \rightarrow \Lambda_c \ell \nu)} \quad (\ell = e, \mu)$$

- longitudinal D^* polarisation

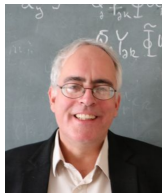
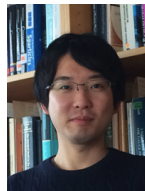
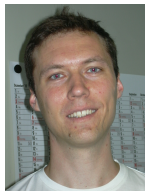
$$F_L(D^*) = \frac{\Gamma(B \rightarrow D_L^* \tau \nu)}{\Gamma(B \rightarrow D^* \tau \nu)} \quad \begin{array}{l} \text{Belle : } 0.60 \pm 0.08 \pm 0.035 \\ \text{SM : } 0.46 \pm 0.04 \end{array}$$

- τ polarisation asymmetries

$$P_\tau(D^{(*)}) = \frac{\Gamma(B \rightarrow D^{(*)} \tau^{\lambda=+1/2} \nu) - \Gamma(B \rightarrow D^{(*)} \tau^{\lambda=-1/2} \nu)}{\Gamma(B \rightarrow D^{(*)} \tau \nu)}$$

- $\text{BR}(B_c \rightarrow \tau \nu)$ – particularly sensitive to scalar contributions

The Crew



MB, CRIVELLIN, DE BOER, KITAHARA, MOSCATI, NIERSTE, NIŠANDŽIĆ
ARXIV:1811.09603

A Closer Look at $B_c \rightarrow \tau\nu$

Constraints on $\text{BR}(B_c \rightarrow \tau\nu)$ advocated in the literature:

- measured total B_c lifetime \triangleright **$\text{BR}(B_c \rightarrow \tau\nu) < 30\%$**

ALONSO, GRINSTEIN, MARTIN CAMALICH (2016)

caveats of τ_{B_c} theory prediction

BENEKE, BUCHALLA (1996)

- large m_c dependence (LO QCD calculation, $1.4 \text{ GeV} < m_c < 1.6 \text{ GeV}$)
- based on heavy quark expansion and non-rel. QCD, but B_c decays dominantly through charm decay

A Closer Look at $B_c \rightarrow \tau\nu$

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ALONSO, GRINSTEIN, MARTIN CAMALICH (2016)

- searches for $B_{u,c} \rightarrow \tau\nu$ at LEP1 $\triangleright \text{BR}(B_c \rightarrow \tau\nu) < 10\%$

AKERROYD, CHEN (2017)

caveats of theory interpretation

- relies crucially on ratio of $b \rightarrow B_c$ vs. $b \rightarrow B_u$ fragmentation functions
- Tevatron and LHC determinations of f_c/f_u not applicable to LEP (hadron collisions vs. Z peak observables)

A Closer Look at $B_c \rightarrow \tau\nu$

Constraints on $\text{BR}(B_c \rightarrow \tau\nu)$ advocated in the literature:

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- searches for $B_{u,c} \rightarrow \tau\nu$ at LEP1 \triangleright **$\text{BR}(B_c \rightarrow \tau\nu) < 10\%$**
AKERROYD, CHEN (2017)

Critical assessment:

- more refined studies needed
- our conservative bound: **$\text{BR}(B_c \rightarrow \tau\nu) < 60\%$**

Effective Hamiltonian

New Physics above B meson scale described model-independently by

$$\mathcal{H}_{\text{eff}}^{\text{NP}} = 2\sqrt{2}G_F V_{cb} \left[(1 + C_V^L) O_V^L + C_S^R O_S^R + C_S^L O_S^L + C_T O_T \right]$$

with the vector, scalar and tensor operators

$$O_V^L = (\bar{c}\gamma^\mu P_L b) (\bar{\tau}\gamma_\mu P_L \nu_\tau)$$

$$O_S^R = (\bar{c}P_R b) (\bar{\tau}P_L \nu_\tau)$$

$$O_S^L = (\bar{c}P_L b) (\bar{\tau}P_L \nu_\tau)$$

$$O_T = (\bar{c}\sigma^{\mu\nu} P_L b) (\bar{\tau}\sigma_{\mu\nu} P_L \nu_\tau)$$

Note: $(\bar{c}\gamma^\mu P_R b) (\bar{\tau}\gamma_\mu P_L \nu_\tau)$ not generated at dimension-six level in the $SU(2)_L \times U(1)_Y$ -invariant theory

Few Technical Remarks

- assume NP only in τ channel – e and μ channels are SM like
- no light right-handed neutrinos
- fit includes $\mathcal{R}(D)$, $\mathcal{R}(D^*)$, $P_\tau(D^*)$, $F_L(D^*)$
- fit uses central values of form factors
 - $B \rightarrow D$ vector and scalar form factors from FLAG WORKING GROUP
 - $B \rightarrow D^*$: V , A_1 , A_2 fit results from HFLAV
 A_0 from BERNLOCHNER ET AL (2017)
 - tensor form factors from BERNLOCHNER ET AL (2017)
 - full set of baryonic $\Lambda_b \rightarrow \Lambda_c$ form factors from DETMOLD ET AL. (2015);
DATTA ET AL. (2017)
- values of Wilson coefficients correspond to scale $\mu = 1$ TeV

One-Dimensional Scenarios

single particle scenarios

C_V^L

left-handed W' boson

left-handed $SU(2)_L$ -singlet vector leptoquark (LQ)

scalar $SU(2)_L$ -triplet and/or -singlet LQ (LH couplings only)

C_S^R

charged Higgs (2HDM-II at large $\tan\beta$)

$SU(2)_L$ -doublet vector LQ

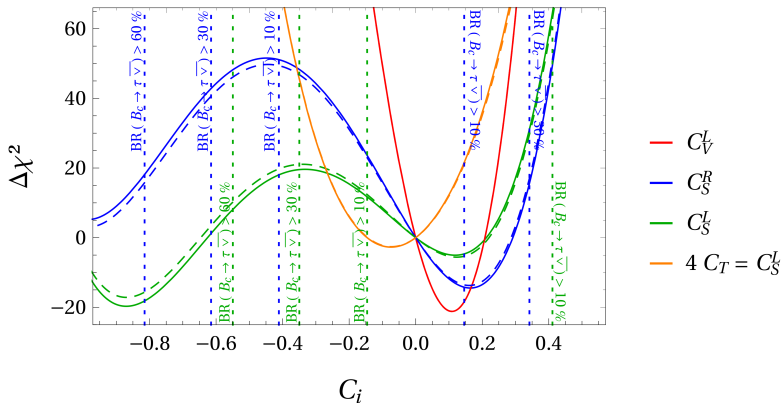
C_S^L

charged Higgs with generic flavor structure

$C_S^L = 4C_T$

scalar $SU(2)_L$ -doublet (relation at NP scale, modified by RG effects)

One-Dimensional Fit Results



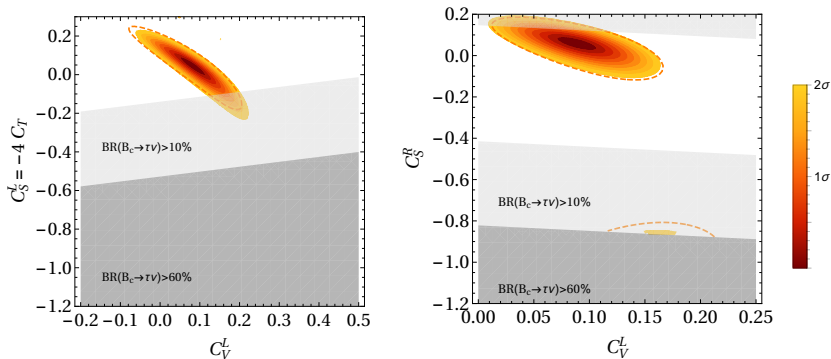
- best fit for $C_V^L \sim 0.11$
- small impact of $F_L(D^*)$ measurement (solid vs. dashed)
- large impact of $\text{BR}(B_c \rightarrow \tau \nu)$ on scalar scenarios

Two-Dimensional Scenarios

single particle scenarios

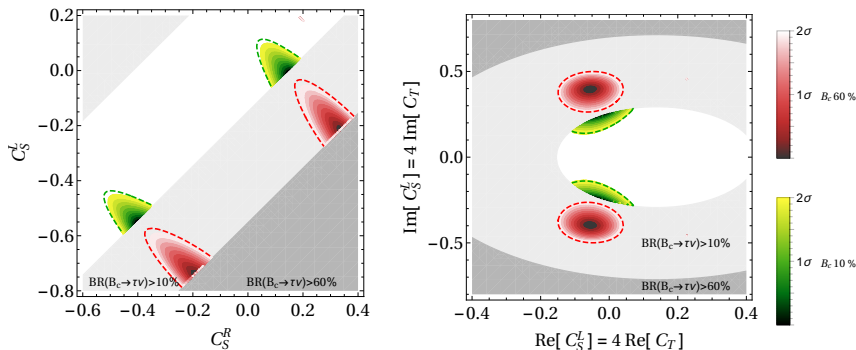
$(C_V^L, C_S^L = -4C_T)$	SU(2) _L -singlet scalar LQ
(C_V^L, C_S^R)	SU(2) _L -singlet vector LQ
(C_S^R, C_S^L)	charged Higgs
$(\text{Re}[C_S^L = 4C_T],$ $\text{Im}[C_S^L = 4C_T])$	scalar SU(2) _L -doublet LQ with CP-violating couplings

Two-Dimensional Fit Results (I)



- good fit for both $(C_V^L, C_S^L = -4C_T)$ and (C_V^L, C_S^R)
- small impact of $BR(B_c \rightarrow \tau\nu)$ constraint

Two-Dimensional Fit Results (II)



- very good fit for (C_S^R, C_S^L) , but only allowed for $BR(B_c \rightarrow \tau\nu) < 60\%$
- good fit for $(C_S^L = 4C_T)$, unless $BR(B_c \rightarrow \tau\nu) < 10\%$ is imposed

The $\Lambda_b \rightarrow \Lambda_c \tau \nu$ Sum Rule

From the phenomenological expressions for $\mathcal{R}(D^{(*)})$ and $\mathcal{R}(\Lambda_c)$, we derive an **approximate sum rule**:

$$\frac{\mathcal{R}(\Lambda_c)}{\mathcal{R}_{\text{SM}}(\Lambda_c)} \simeq 0.262 \frac{\mathcal{R}(D)}{\mathcal{R}_{\text{SM}}(D)} + 0.738 \frac{\mathcal{R}(D^*)}{\mathcal{R}_{\text{SM}}(D^*)} + \mathcal{O}(10^{-2})$$

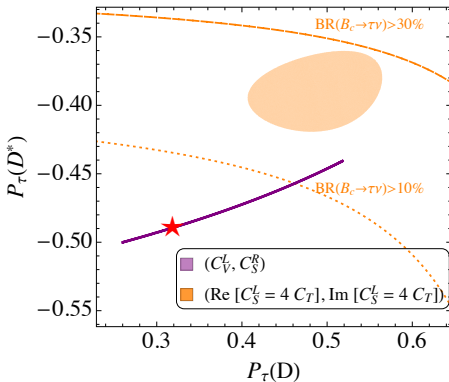
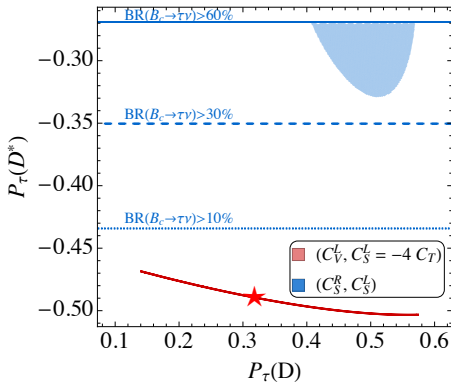
- enhancement of $\mathcal{R}(D^{(*)})$ implies $\mathcal{R}(\Lambda_c) > \mathcal{R}_{\text{SM}}(\Lambda_c) = 0.33 \pm 0.01$
- model-independent prediction from current $\mathcal{R}(D^{(*)})$ data:

$$\mathcal{R}(\Lambda_c) = 0.41 \pm 0.02_{\mathcal{R}(D^{(*)})} \pm 0.01_{\text{form factors}}$$

- **experimental cross-check of $\mathcal{R}(D^{(*)})$ anomaly**

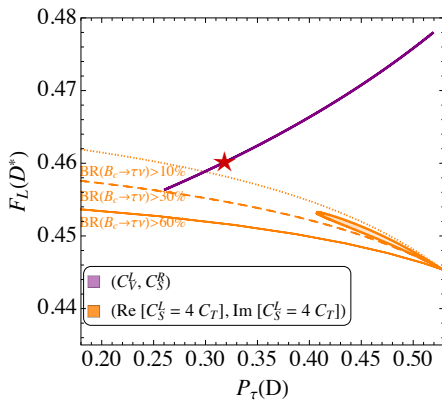
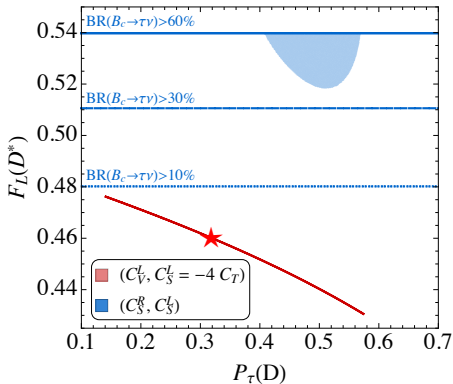
Correlations between Polarization Observables (I)

$P_\tau(D)$ and $P_\tau(D^*)$



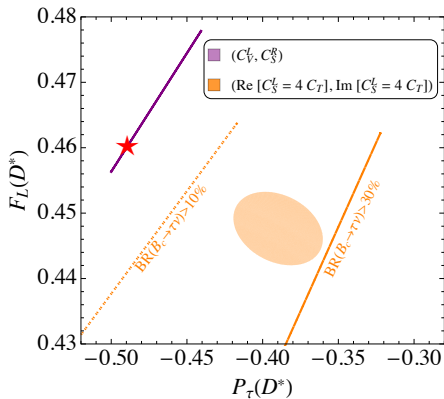
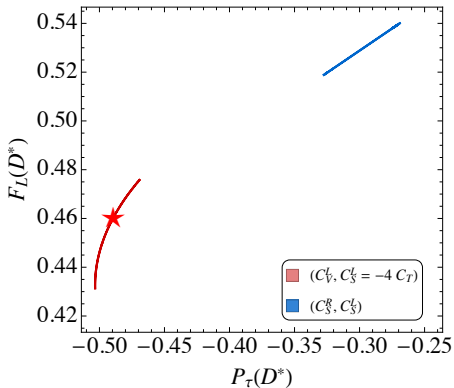
Correlations between Polarization Observables (II)

$P_\tau(D)$ and $F_L(D^*)$



Correlations between Polarization Observables (III)

$P_\tau(D^*)$ vs. $F_L(D^*)$



Summary: What's New in $b \rightarrow c\tau\nu$

- updated 1D and 2D fits, including recent $F_L(D^*)$ measurement
- critical assessment of $B_c \rightarrow \tau\nu$ constraint
- $\Lambda_b \rightarrow \Lambda_c\tau\nu$ provides experimental cross-check of $\mathcal{R}(D^{(*)})$ anomaly
- polarisation observables well suited to distinguish among different EFT scenarios ➤ requires better understanding of form factors in addition to decently precise measurements